

rans(eXtreme)

Analysis framework for multi-fluid compressible hydrodynamic simulations

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Physics motivation

- Limitations of the current 1D modelling of turbulence in stars
- Closures for approximated or neglected physics
 - what do we actually neglect?
- Hydrodynamic stellar structure equations (time-dependent, non-local)
 - (no rotation, no magnetic fields)

Computational motivation

- Comprehensive analysis of hydrodynamic simulations done at runtime and user-friendly post-processing

Structure

- **Theory:** Reynolds and Favrian decomposition

$$A(r, \theta, \phi) = \overline{A}(r) + A'(r, \theta, \phi) \quad \overline{A}(r) = \frac{1}{\Delta T \Delta \Omega} \int_{\Delta T} \int_{\Delta \Omega} A(r, \theta, \phi) dt d\Omega$$

$$F(r, \theta, \phi) = \tilde{F}(r) + F''(r, \theta, \phi) \quad \tilde{F} = \overline{\rho F} / \bar{\rho}$$

$$\underbrace{\overline{u}_r}_{\text{mean velocity}} = \underbrace{\tilde{u}_r}_{\text{expansion velocity } \partial_t M / 4\pi r^2 \bar{\rho}} - \underbrace{\overline{u''_r}}_{\text{turbulent mass flux } -\overline{\rho' u'_r} / \bar{\rho}}$$

- <https://github.com/mmicromegas/ransX/tree/master/DOCS>
- **Hydrodynamic Code Implementation:** calculation of mean-fields at runtime of simulation
 - https://github.com/mmicromegas/ransX/tree/master/UTILS/FOR_YOUR_HYDRO
- **Post-Processing in Python**
 - <https://github.com/mmicromegas/ransX>

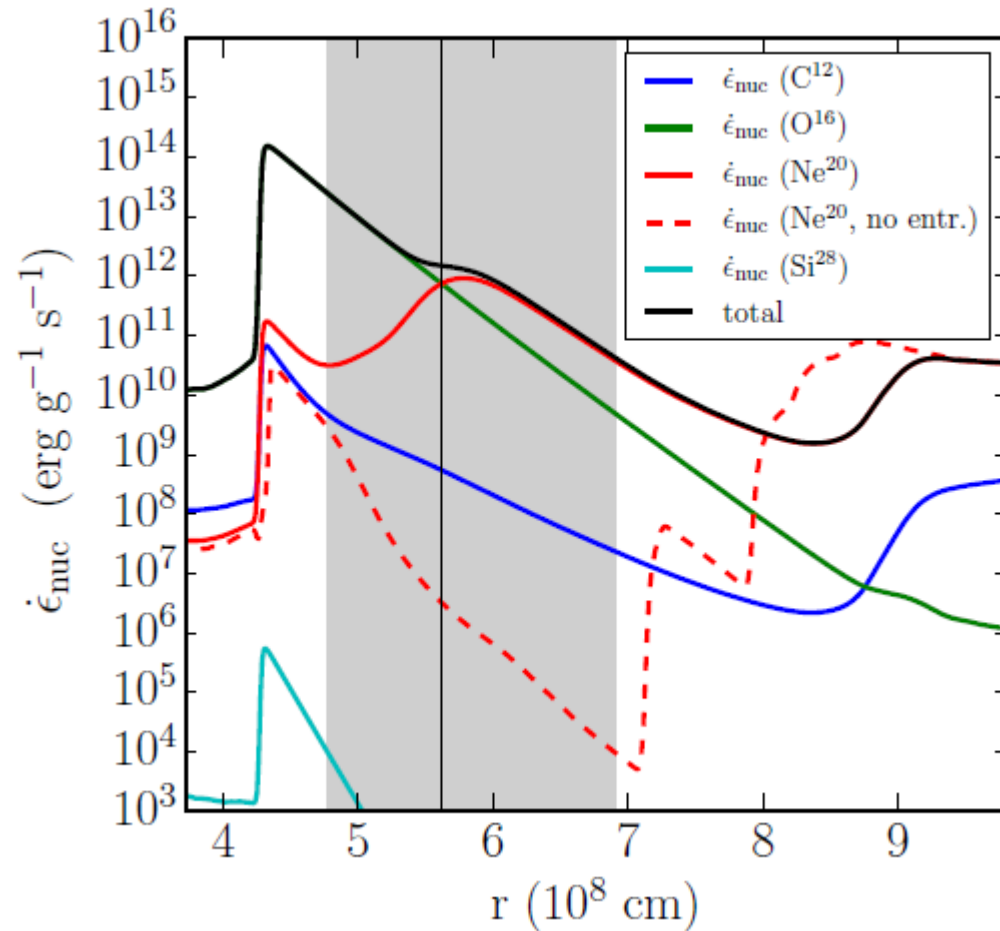
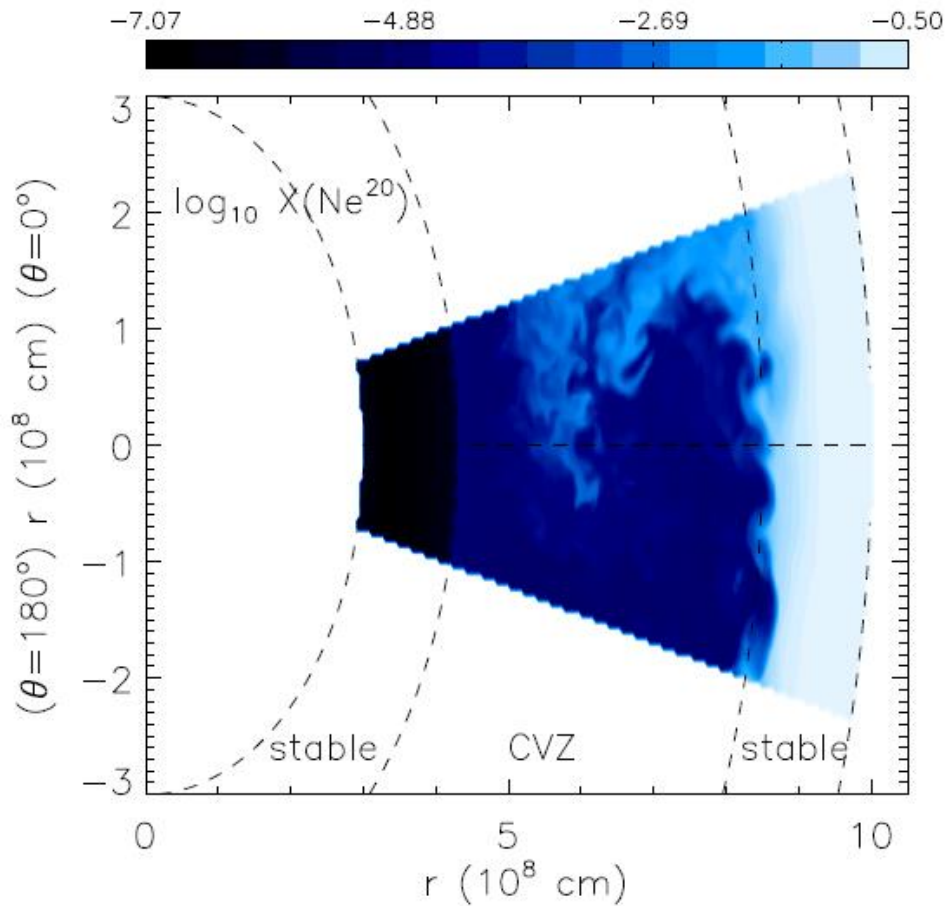
Results

- Transport/Flux/Variance equations for evolution of mass, momenta, kinetic/internal/total energy, temperature, enthalpy, pressure
 - Transport/Flux/Variance equations for evolution of chemical composition
 - Eulerian diffusivities (to guide us towards new composition mixing models)
 - Hydrodynamic stellar structure equations (3 versions)
 - general
 - simplified (based on flux evolution equations)
 - simplified (for adiabatic flow in HSE)
- * all of them well validated with our oxygen-neon burning simulation
- * for more details see <https://github.com/mmicromegas/ransX/tree/master/DOCS/RANDOM>

Oxygen-Neon burning convective shell

[2018MNRAS.481.2918M](#) Mocák et al, 2018

- multiple burning zones within single convection zone



Hydrodynamic stellar structure equations

- Below is a complete set of hydrodynamic stellar structure equations derived from RANS equations (viscosity explicitly neglected), where red terms are the ones used in classical approach:

$$\partial_r \bar{m} = 4\pi r^2 \bar{\rho} + (4\pi r^3 / 3 \tilde{u}_r) [-\nabla_r f_\rho + (f_\rho / \bar{\rho}) \partial_r \bar{\rho} - \bar{\rho} \bar{d} - \partial_t \bar{\rho}]$$

$$\partial_r \bar{P} = \bar{\rho} \tilde{g} - \bar{\rho} \partial_t \tilde{u}_r - \nabla_r \tilde{R}_{rr} - \bar{G}_r^M - \bar{\rho} \tilde{u}_r \partial_r \tilde{u}_r$$

$$\partial_r \tilde{L} = 4\pi r^2 \tilde{\rho} \tilde{\epsilon}_{nuc} + 4\pi r^2 \left[-\nabla_r (f_i + f_{th} + f_K + f_p) - \bar{P} \bar{d} - \tilde{R}_{ir} \partial_r \tilde{u}_i + W_b + \bar{\rho} \tilde{D}_t \tilde{u}_i \tilde{u}_i / 2 - \bar{\rho} \partial_t \tilde{\epsilon}_t \right] + \tilde{\epsilon}_t \partial_r 4\pi r^2 \tilde{\rho} \tilde{u}_r$$

$$\partial_r \bar{T} = (1 / \bar{u}_r) [-\nabla_r f_T + (1 - \Gamma_3) \bar{T} \bar{d} + (2 - \Gamma_3) \bar{T}' \bar{d}' + \epsilon_{nuc} / c_v + \nabla \cdot f_{th} / (\rho c_v) - \partial_t T]$$

$$\partial_t \tilde{X}_i = \tilde{X}_i^{nuc} - (1 / \bar{\rho}) \nabla_r f_i - \tilde{u}_r \partial_r \tilde{X}_i$$

<https://github.com/mmicromegas/ransX/blob/master/DOCS/RANDOM/hsse.pdf>

https://github.com/mmicromegas/ransX/blob/master/DOCS/RANDOM/hsse_alternative.pdf

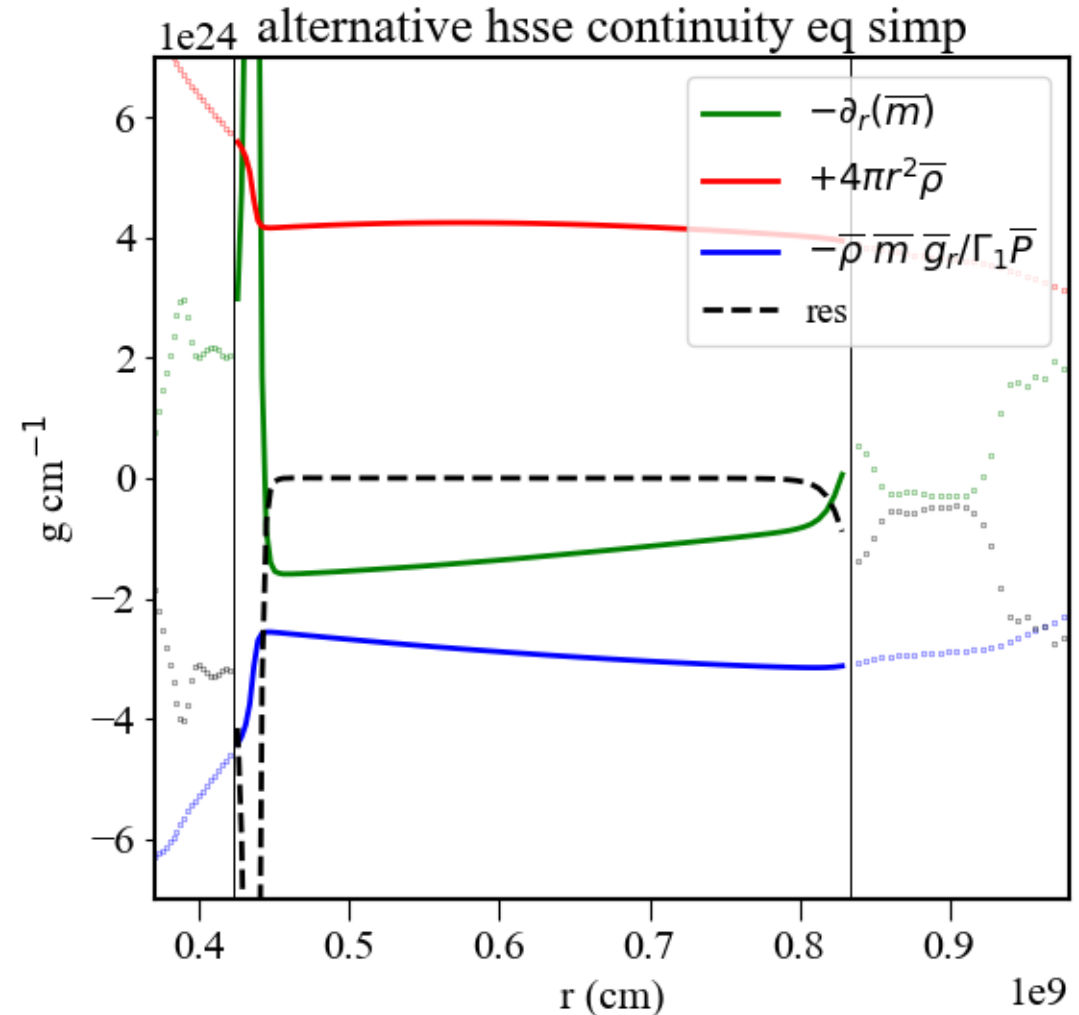
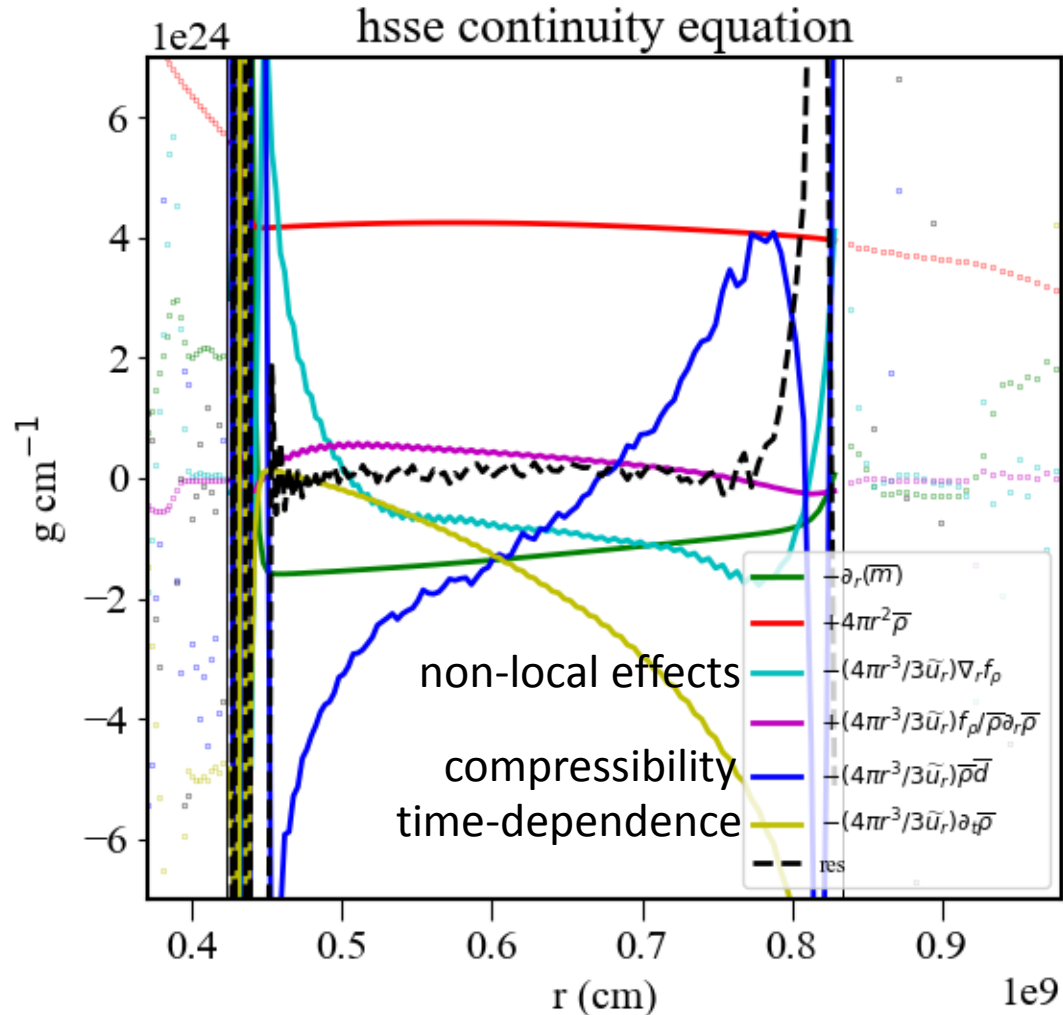
Hydrodynamic stellar structure equations

- Stellar gradients and dilatation flux $\tilde{R}_{rr} \partial_r \bar{Q} \sim -\bar{\rho} \bar{Q} \overline{u'_r d''}$ (inferred from flux equations)
- Below is a set of alternative hydrodynamic stellar structure equations derived from the relation between stellar gradients and dilatation flux, where the Q was replaced by density (ρ), pressure (P), total energy (e_t) and temperature (T) [composition X equation is standard continuity equation]

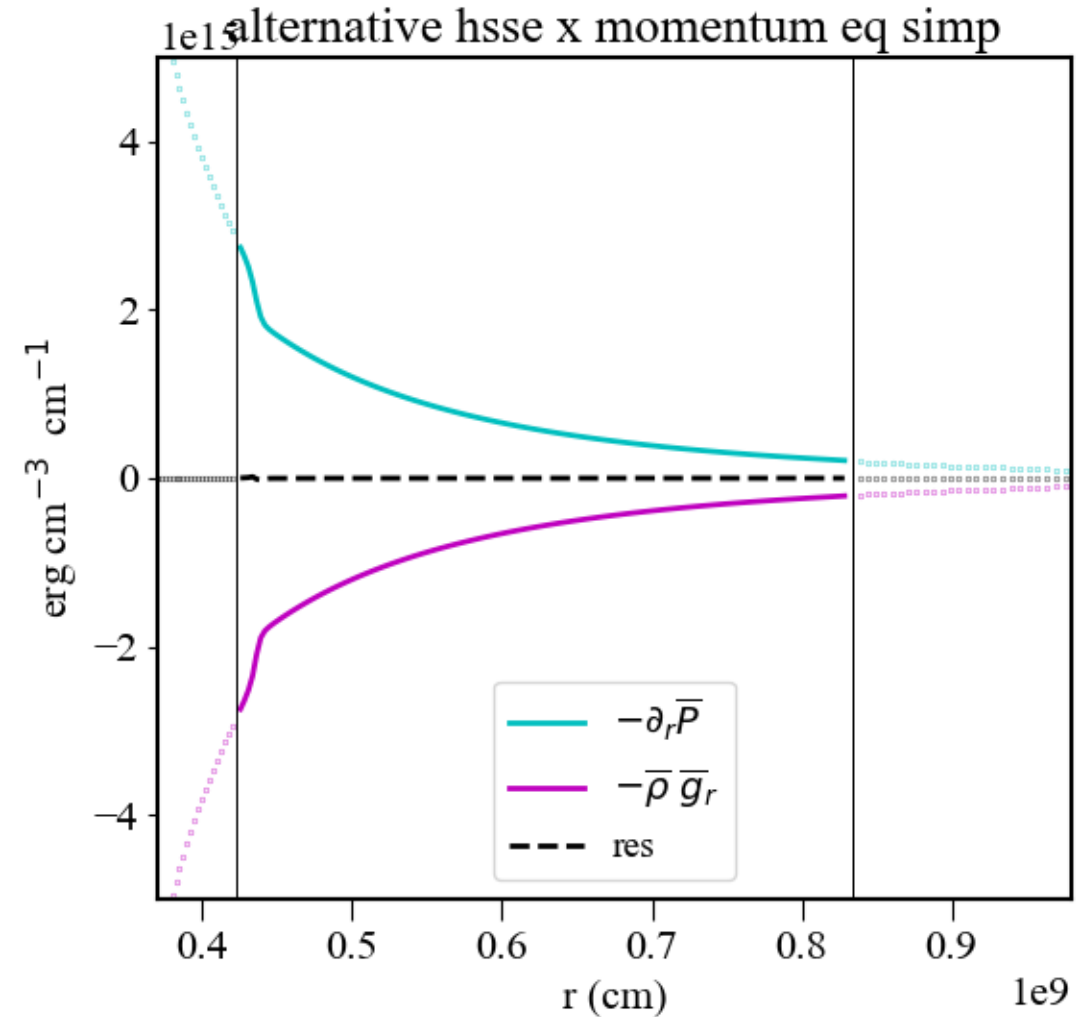
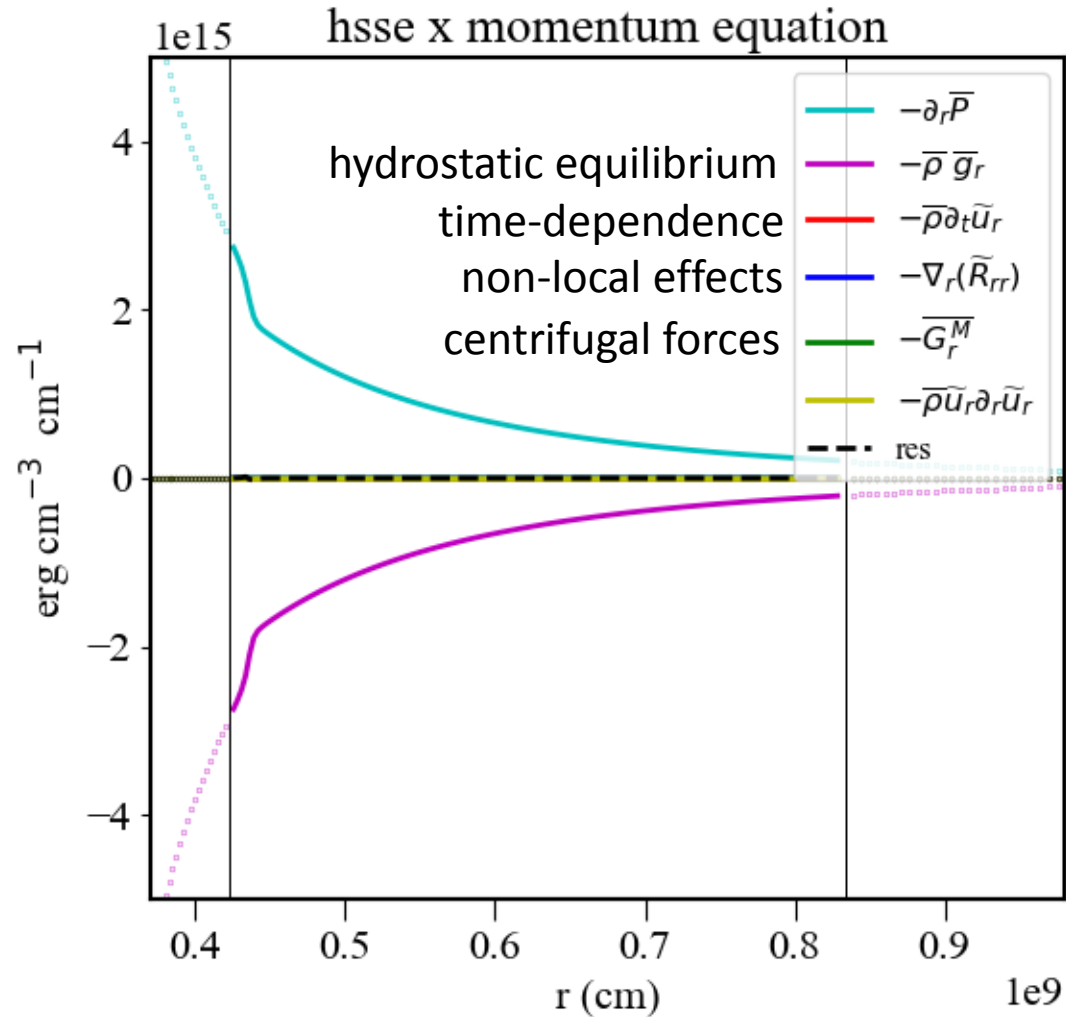
$$\begin{array}{l}
 \partial_r \bar{m} = -\bar{\rho} \bar{m} \overline{u'_r d''} / \tilde{R}_{rr} + 4\pi r^2 \bar{\rho} \\
 \partial_r \bar{P} = -\Gamma_1 \bar{\rho} \bar{P} \overline{u'_r d''} / \tilde{R}_{rr} \\
 \partial_r \tilde{L} = +\tilde{\epsilon}_t \partial_r 4\pi r^2 \bar{\rho} \tilde{u}_r - 4\pi r^2 \bar{\rho} \tilde{u}_r \bar{P} \overline{u'_r d''} / \tilde{R}_{rr} \\
 \partial_r \bar{T} = -(\Gamma_3 - 1) \bar{\rho} \bar{T} \overline{u'_r d''} / \tilde{R}_{rr} \\
 \partial_t \tilde{X}_i = \tilde{X}_i^{nuc} - (1/\bar{\rho}) \nabla_r f_i - \tilde{u}_r \partial_r \tilde{X}_i \\
 \tilde{u}_r = -\partial_t \bar{M} / 4\pi r^2 \bar{\rho}
 \end{array}
 \xrightarrow[\overline{u'_r d''} \sim \frac{\tilde{R}_{rr} \bar{g}_r}{\Gamma_1 \bar{P}}]{\partial_r \bar{P} \sim -\bar{\rho} \bar{g}_r}
 \begin{array}{l}
 \partial_r \bar{m} = -\bar{\rho} \bar{m} \bar{g}_r / \Gamma_1 \bar{P} + 4\pi r^2 \bar{\rho} \\
 \partial_r \bar{P} = -\bar{\rho} \bar{g}_r \\
 \partial_r \tilde{L} = -4\pi r^2 \tilde{u}_r \bar{\rho} \bar{g}_r / \Gamma_1 + \tilde{\epsilon}_t \partial_r 4\pi r^2 \bar{\rho} \tilde{u}_r \\
 \partial_r \bar{T} = -(\Gamma_3 - 1) \bar{\rho} \bar{T} \bar{g}_r / \Gamma_1 \bar{P} \\
 \partial_t \tilde{X}_i = \tilde{X}_i^{nuc} - (1/\bar{\rho}) \nabla_r f_i - \tilde{u}_r \partial_r \tilde{X}_i
 \end{array}$$

$$m = \rho V = \rho \frac{4}{3} \pi r^3 \quad M = \int \rho(r) dV$$

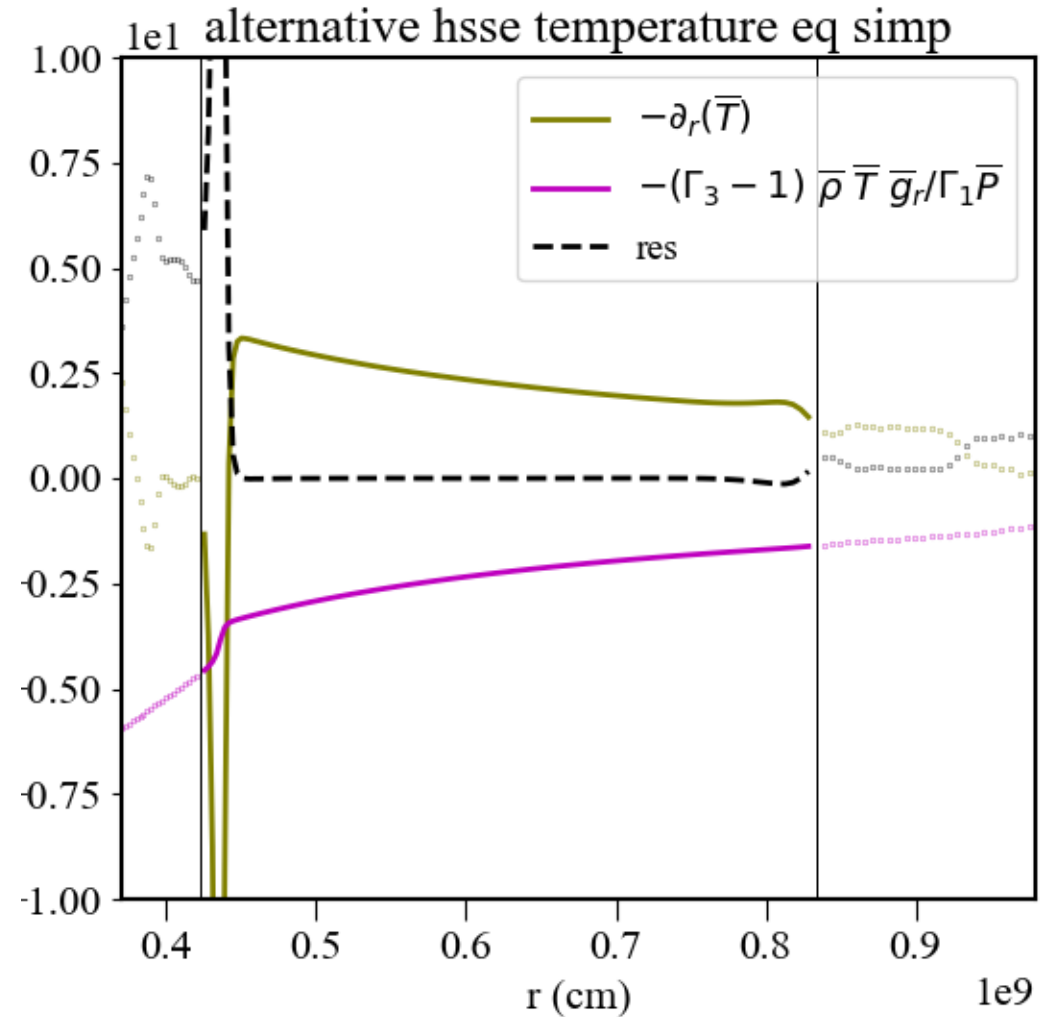
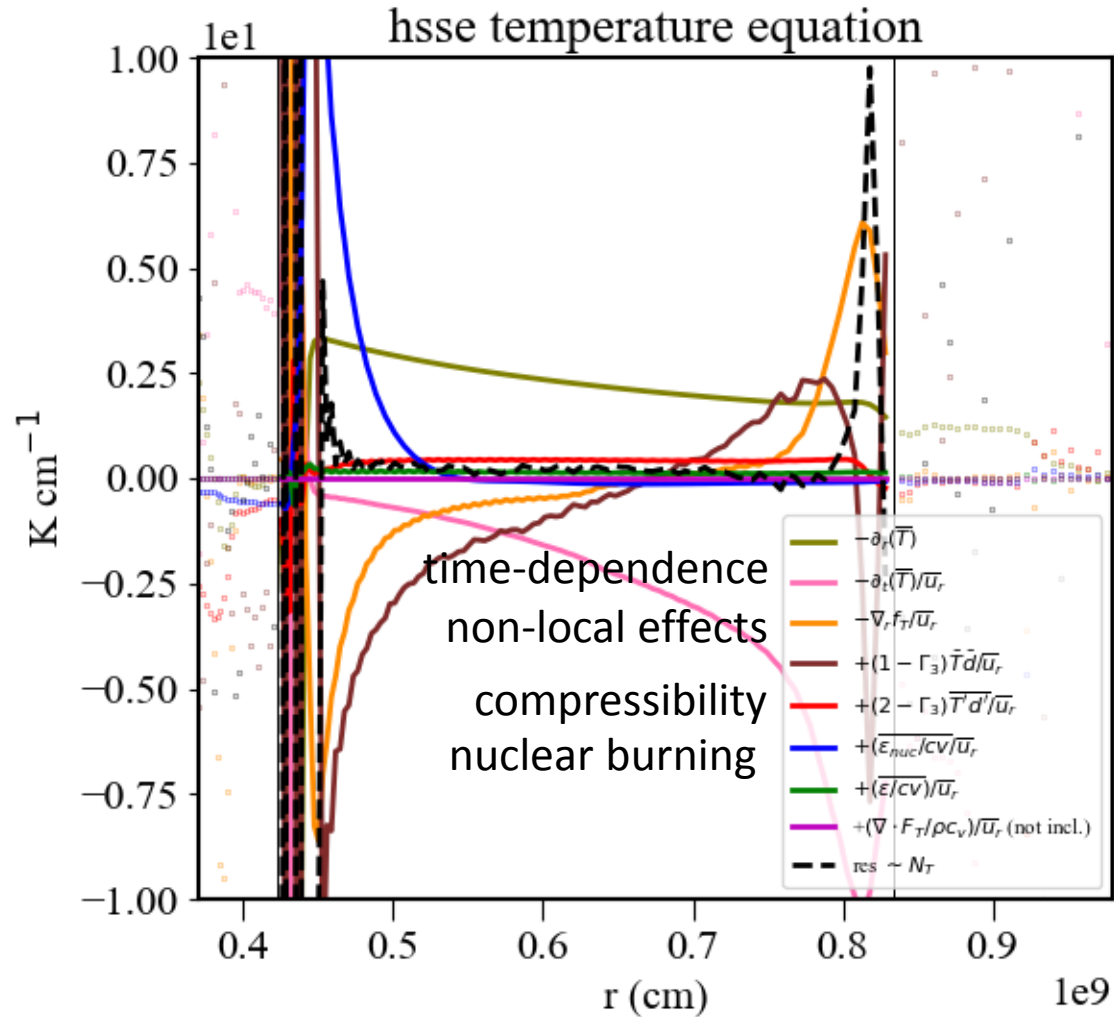
Hydrodynamic stellar structure equations



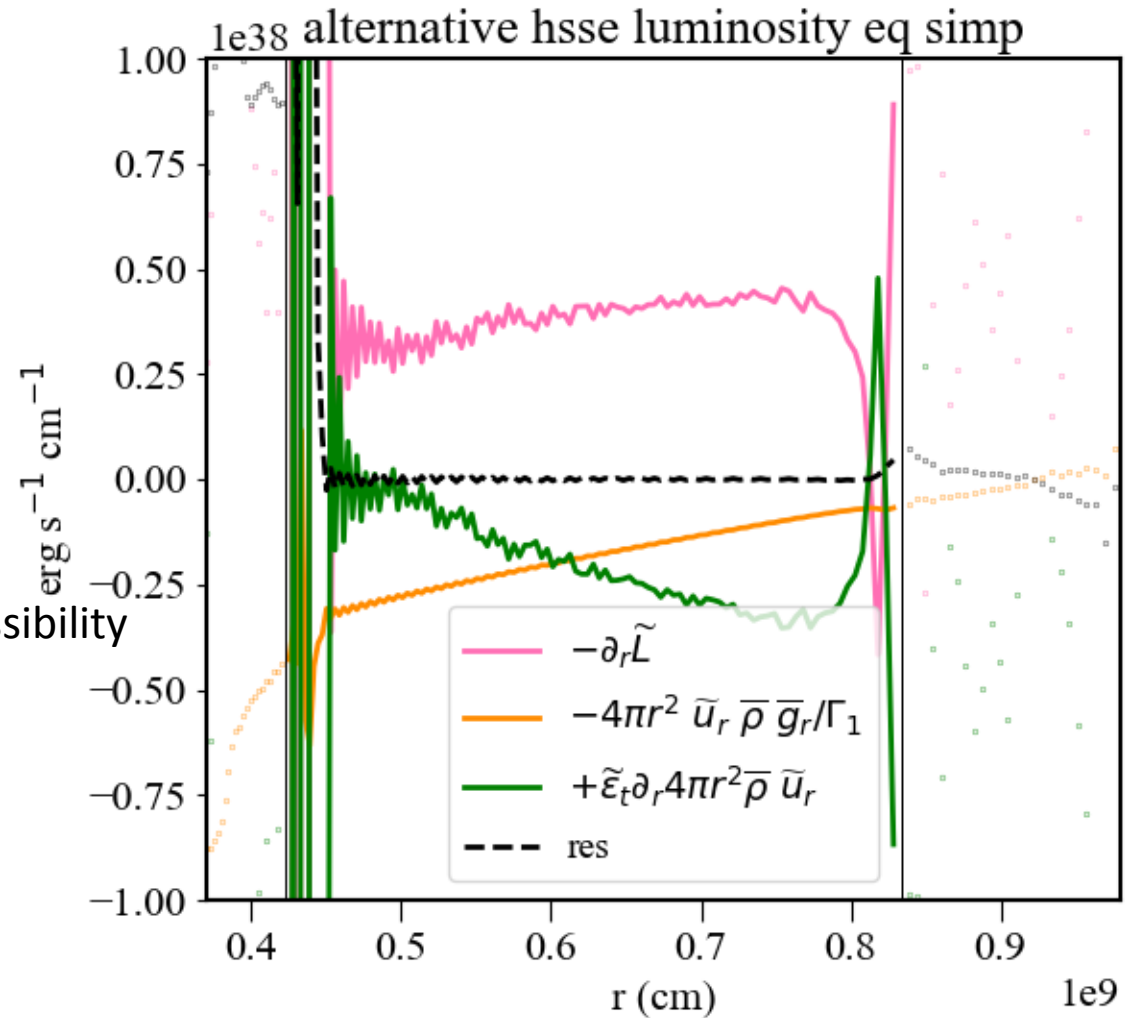
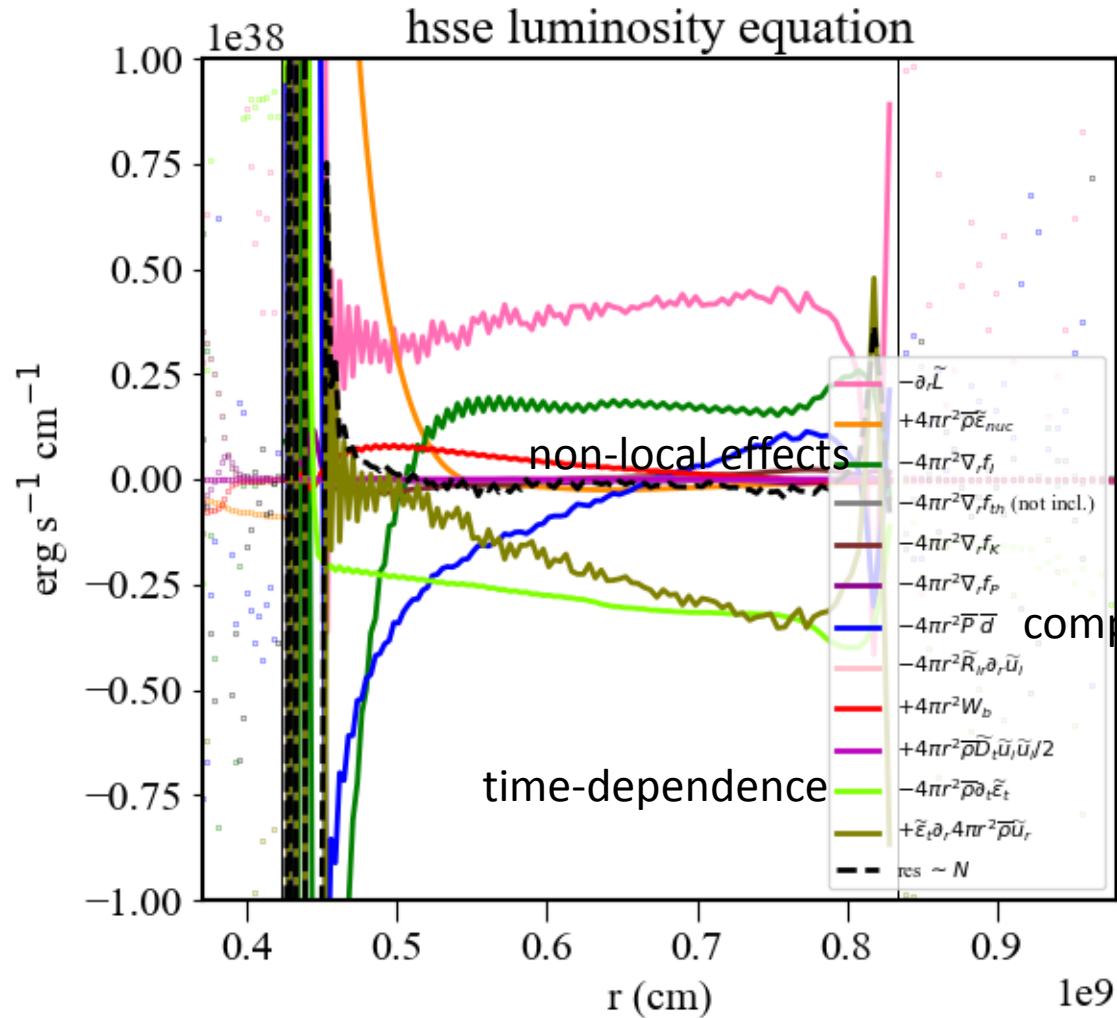
Hydrodynamic stellar structure equations



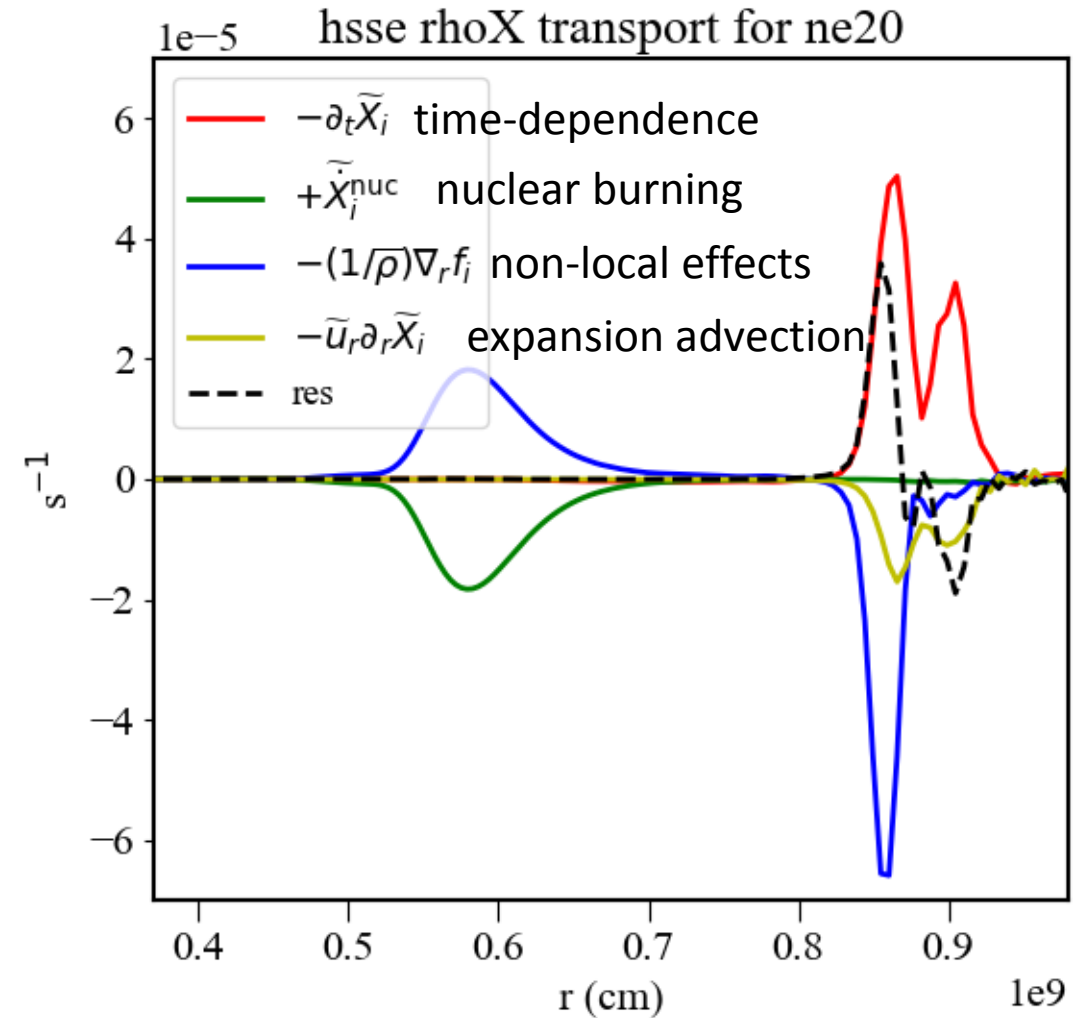
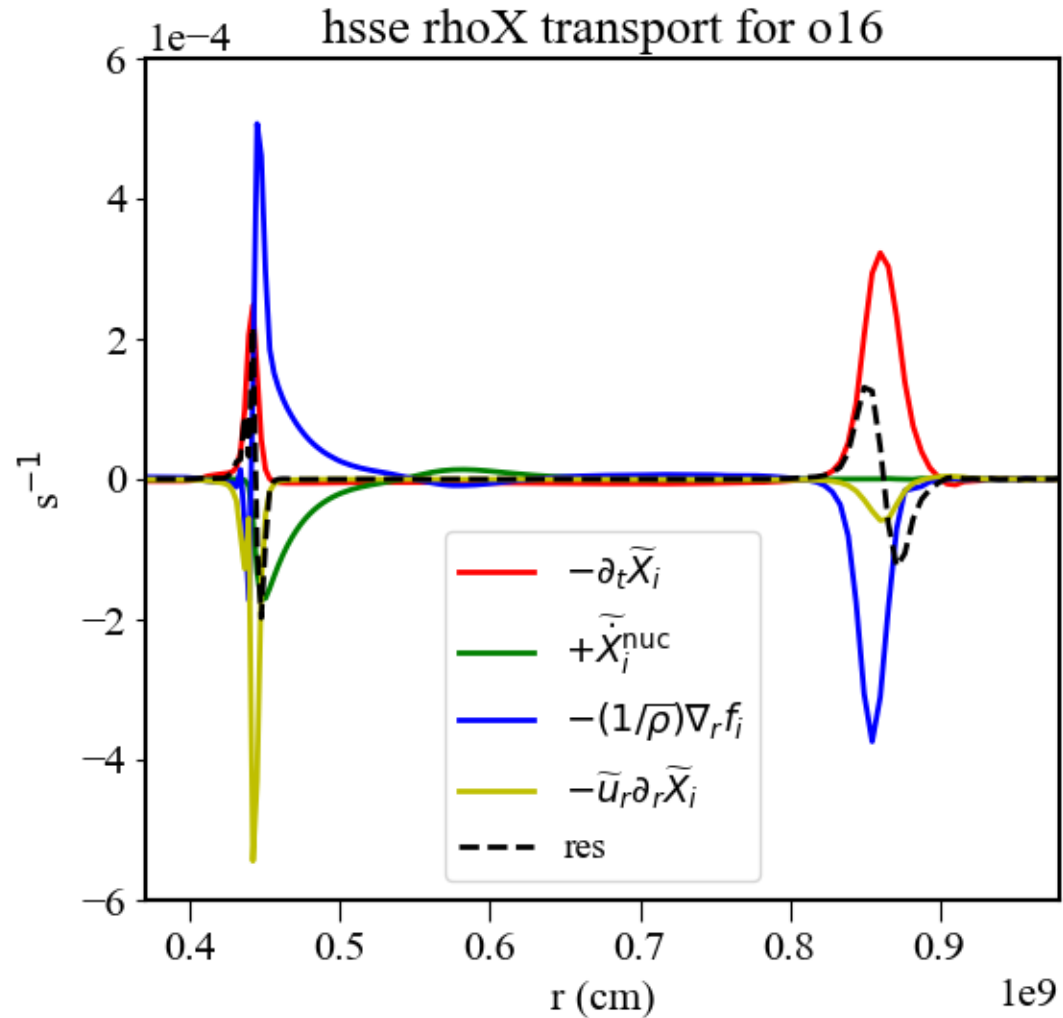
Hydrodynamic stellar structure equations



Hydrodynamic stellar structure equations



Hydrodynamic stellar structure equations



Composition flux model with Gaussian eddy diffusivity

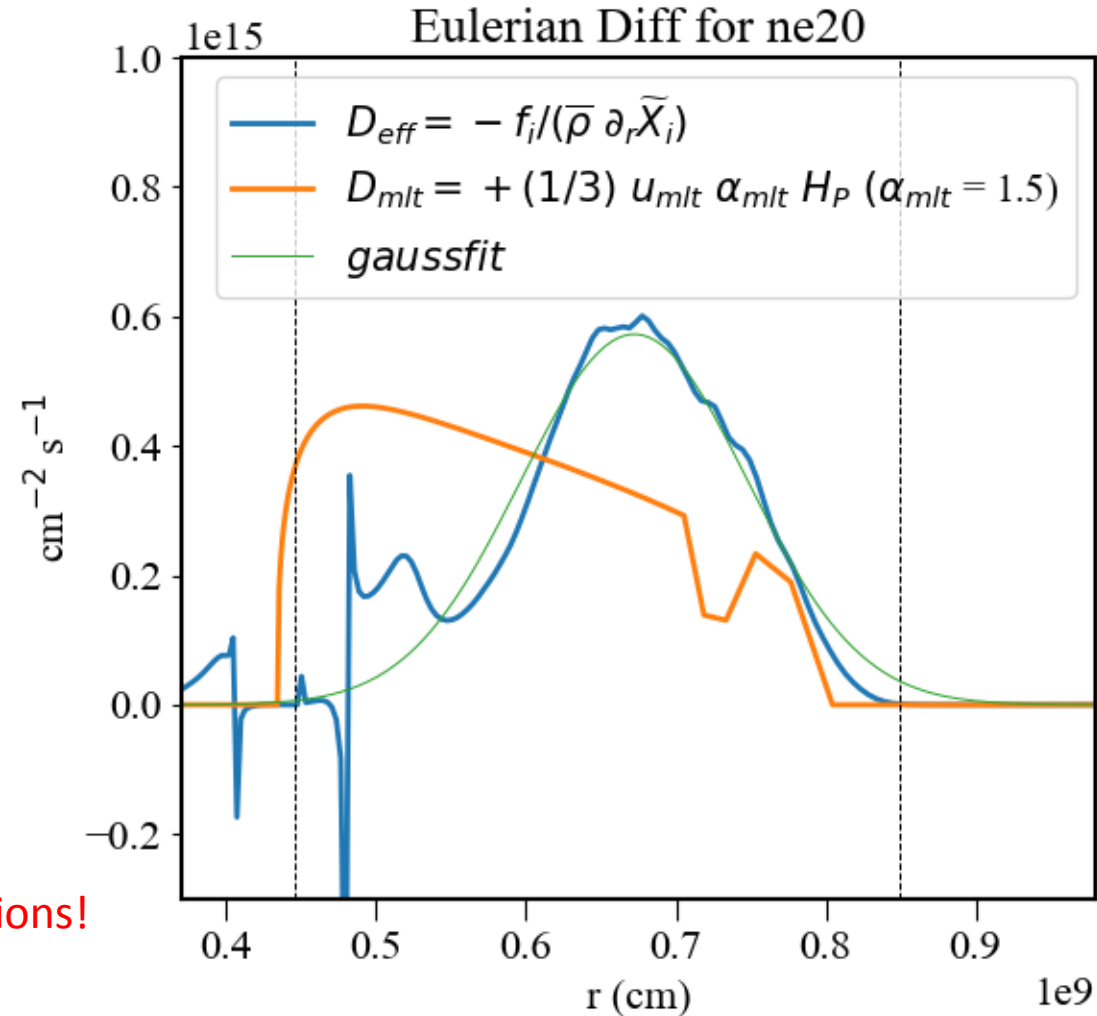
$$f_i = -D \bar{\rho} \partial_r \tilde{X}_i$$

$$D_{mlt} = \frac{1}{3} u_{mlt} (\alpha H_P)$$

$$D_{eff} = -f_i / (\bar{\rho} \partial_r \tilde{X}_i)$$

$$D_{gauss} = \max(D_{mlt}) e^{-\frac{(r-r_c^{middle})^2}{2 \text{width}_c^2}}$$

To get this right is essential, because in reactive flows, **mixing controls rate of nuclear reactions!**



Summary and results

- Analysis framework for multi-fluid compressible hydrodynamic simulations completed <https://github.com/mmicromegas/ransX> (for cartesian and spherical geometry only, no rotation, no magnetic fields)
- Time-dependence, non-locality and compressibility effects play important roles during life of stars
- Transport-diffusion model for composition flux requires Gaussian-like eddy- diffusivity (fluxes of some active elements, where nuclear burning significantly affects their mean gradients are even more complicated) - MLT based diffusivities have limitations
- ransx@googlegroups.com (if interested, please send an email to miroslav.mocak@gmail.com to be added to our group)

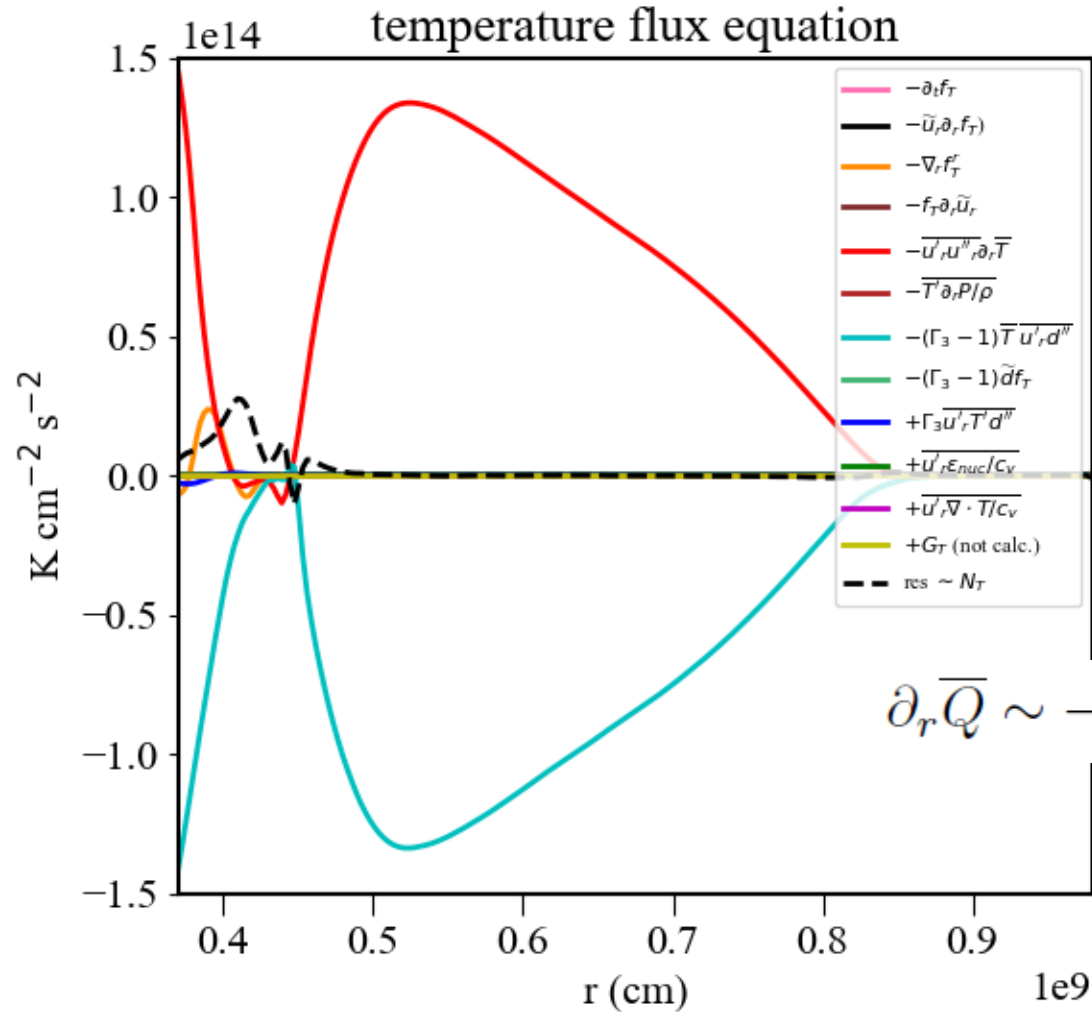
Future plans (to science the shit of out this)

- Look for potential closures of unknowns in general hydrodynamic stellar structure equations in engineering literature and atmospheric sciences for turbulent regions, their boundaries and stable layers
- Help ransX to become standard for all stellar hydrodynamic simulations including core, envelope and atmospheric convection
- Extend library of our ransX hydrodynamic simulations with core helium flash, dual core flash, core carbon flash, O-Ne-C burning with two distinct convection shells – all setups prepared in PROMPI already
- Incorporate the Gaussian eddy-diffusivity mixing model to 1D stellar code (e.g. MESA)

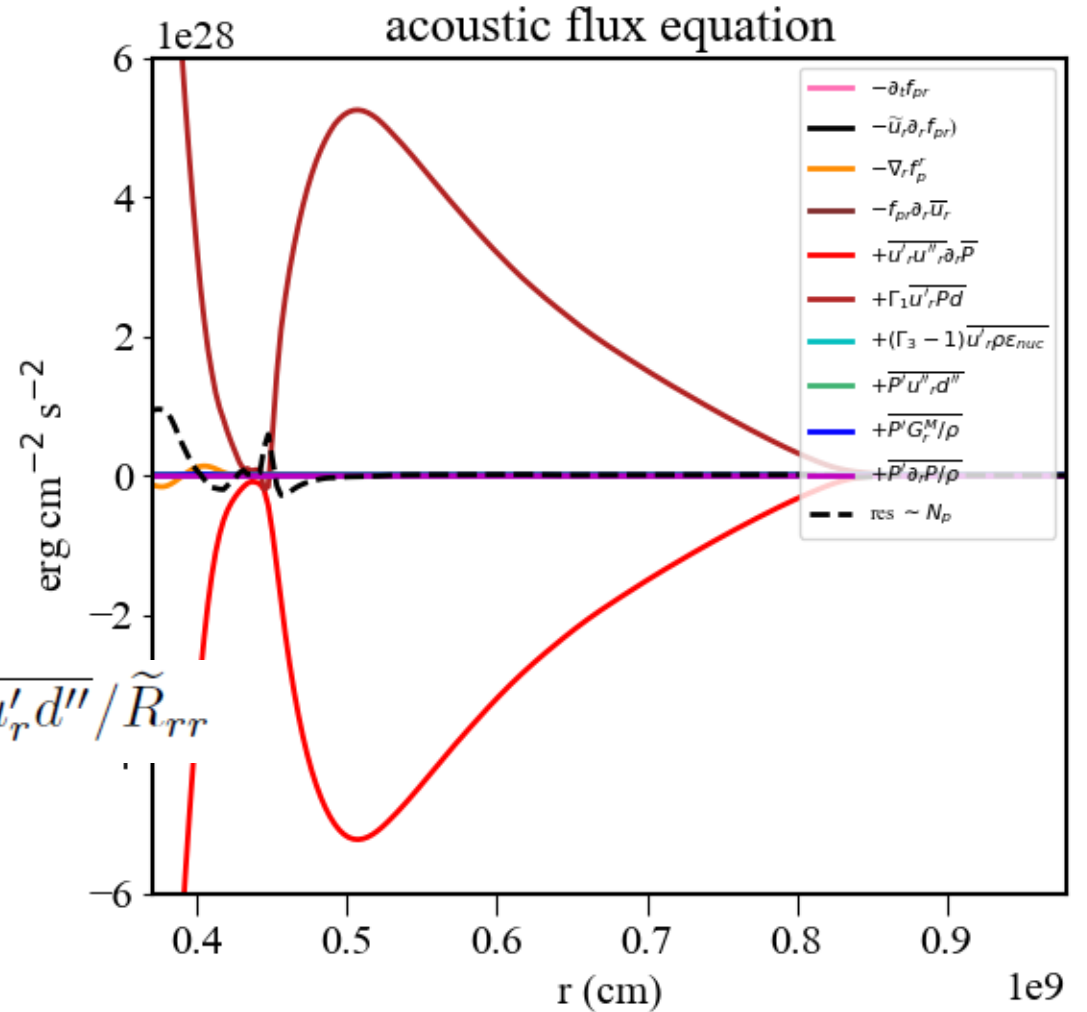
Contact: miroslav.mocak@gmail.com

THANKS

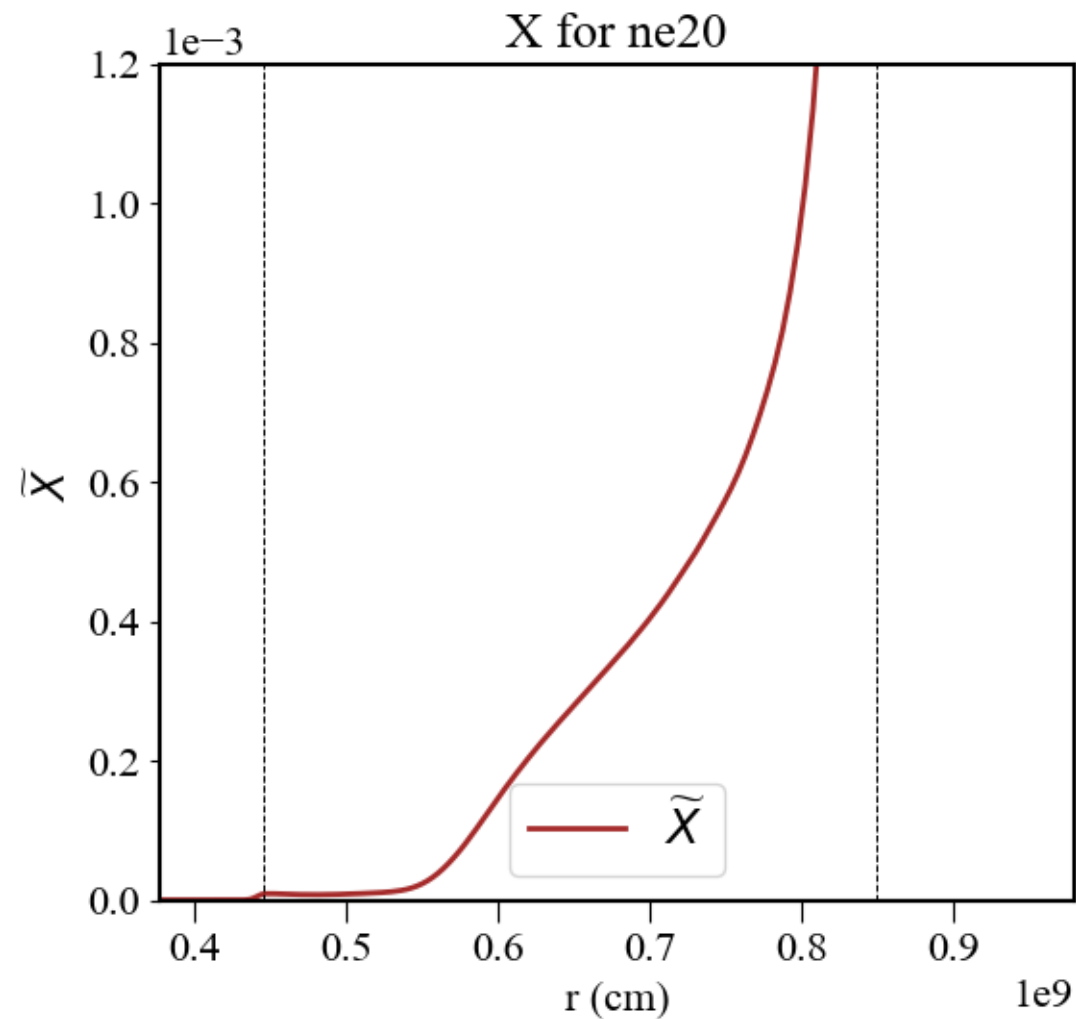
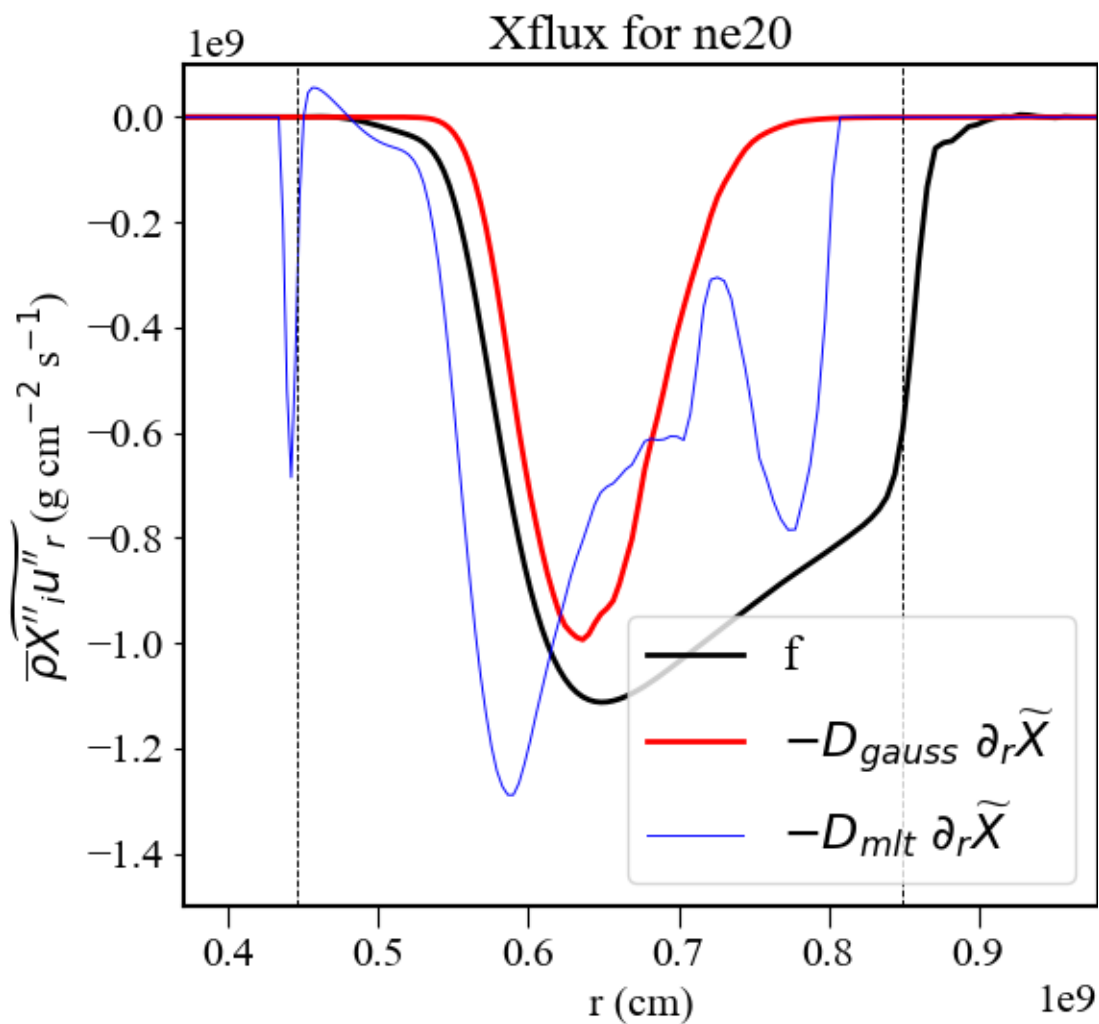
Flux evolution equations and stellar gradients



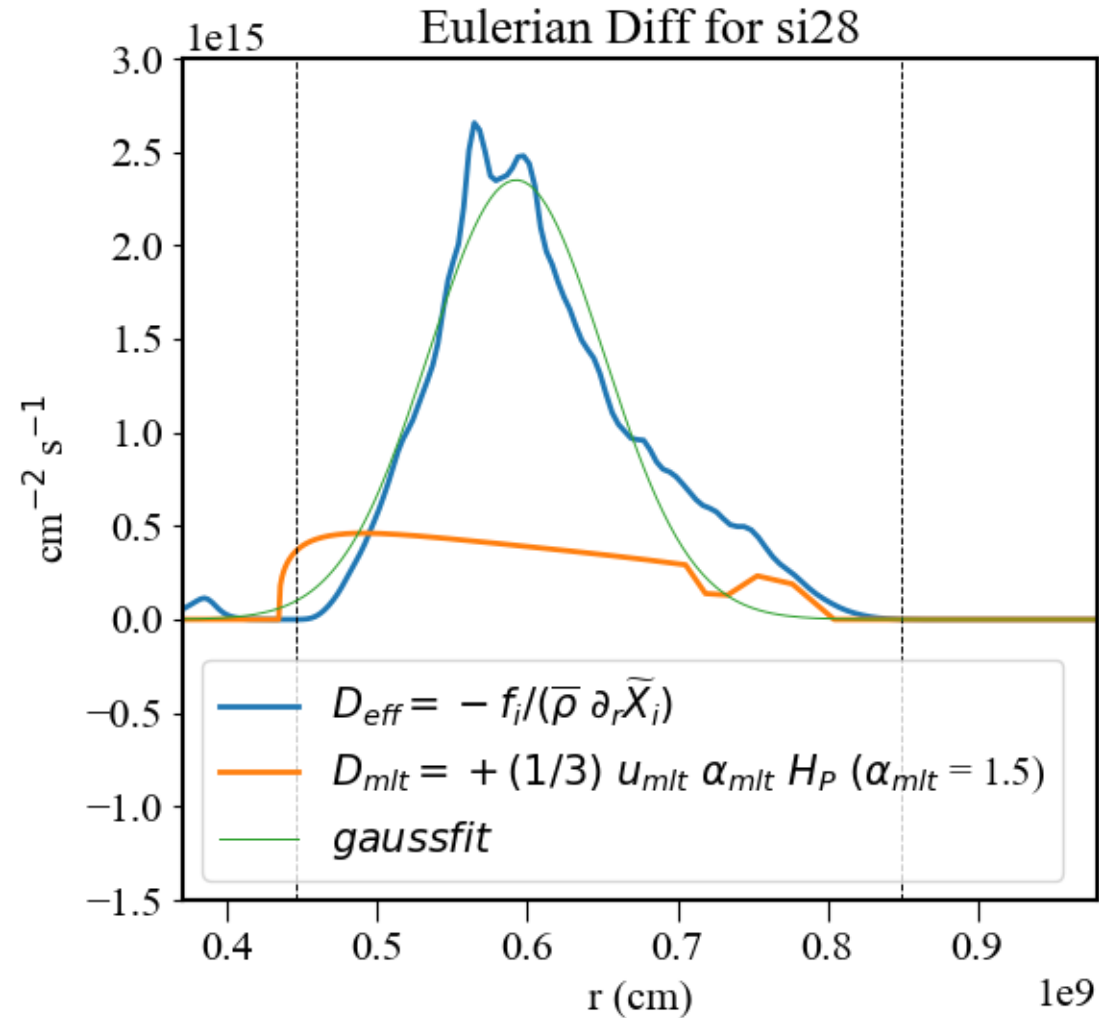
$$\partial_r \bar{Q} \sim -\bar{\rho} \bar{Q} \overline{u'_r d''} / \tilde{R}_{rr}$$



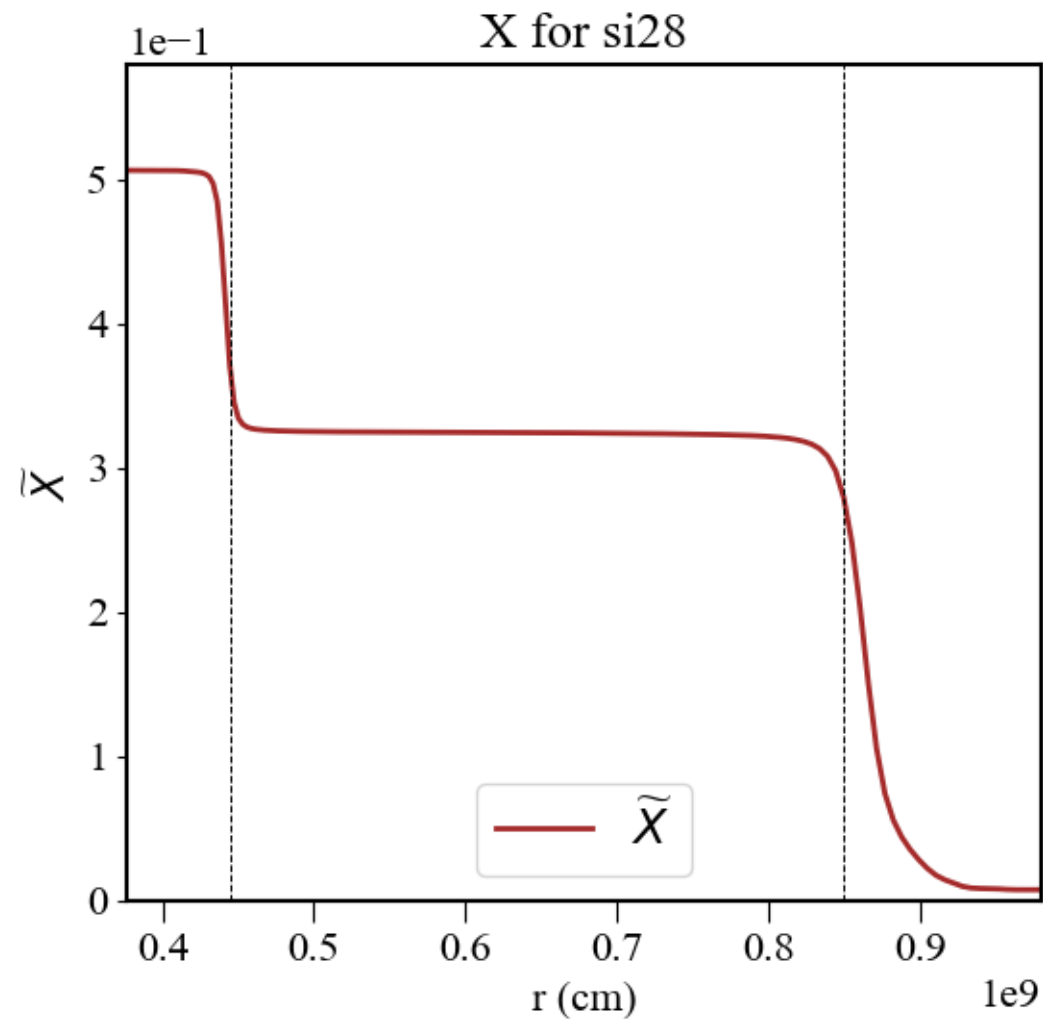
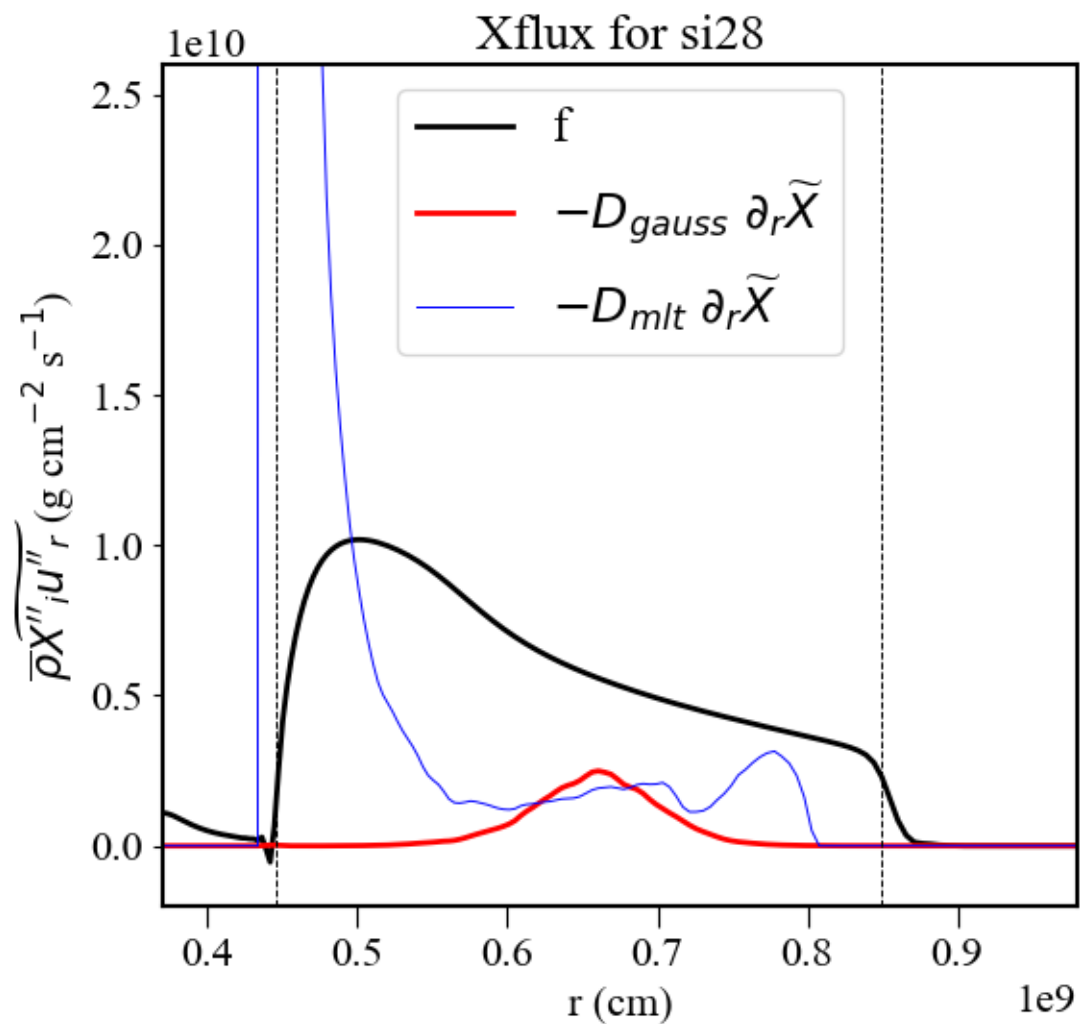
Composition flux model



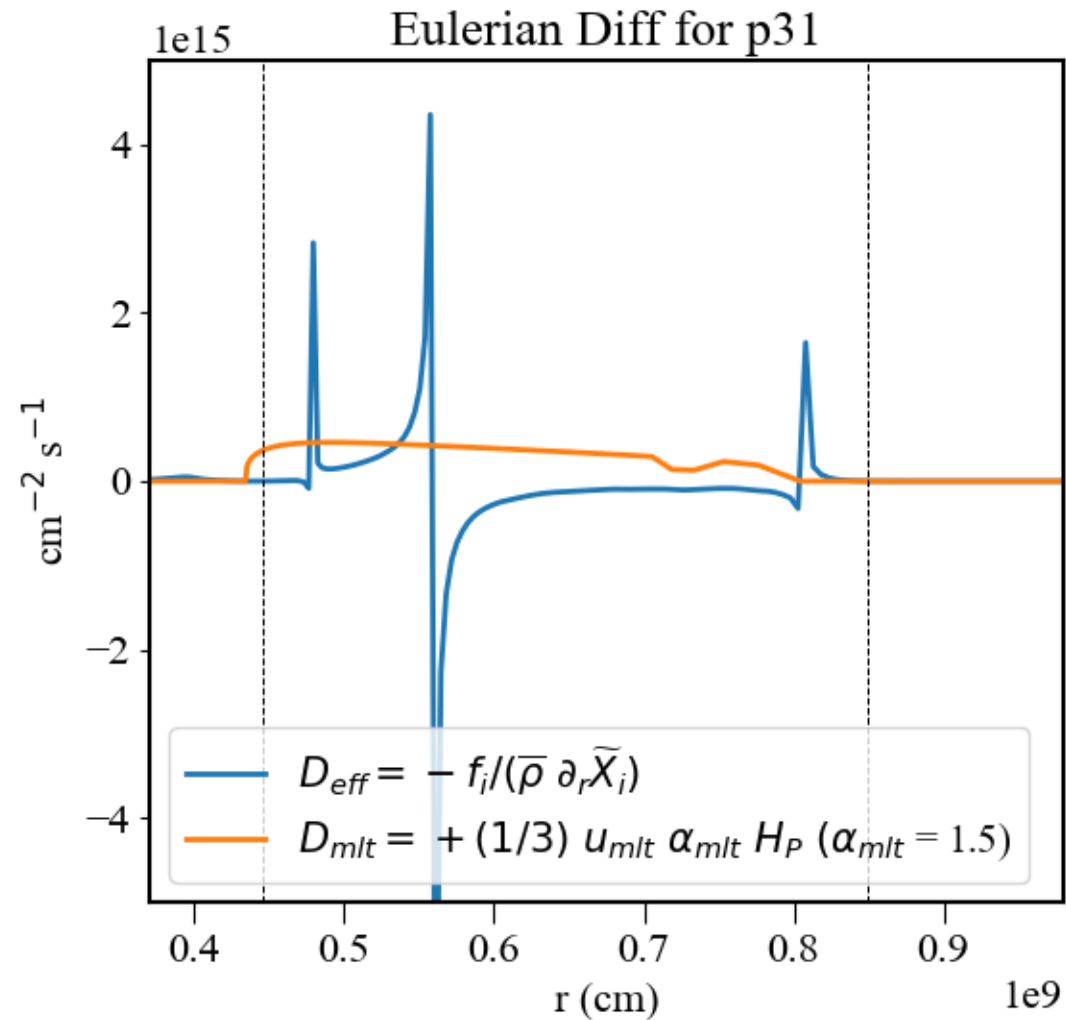
Composition flux model



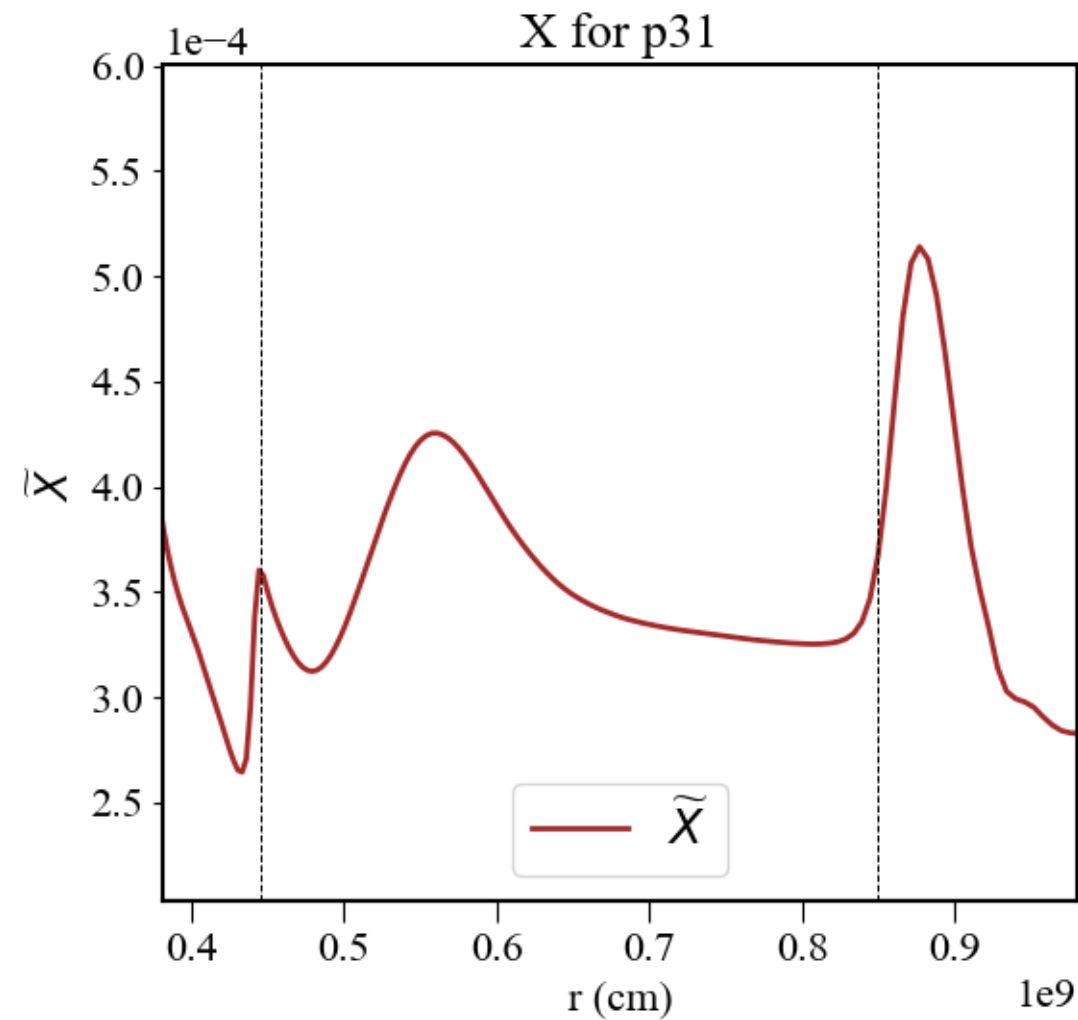
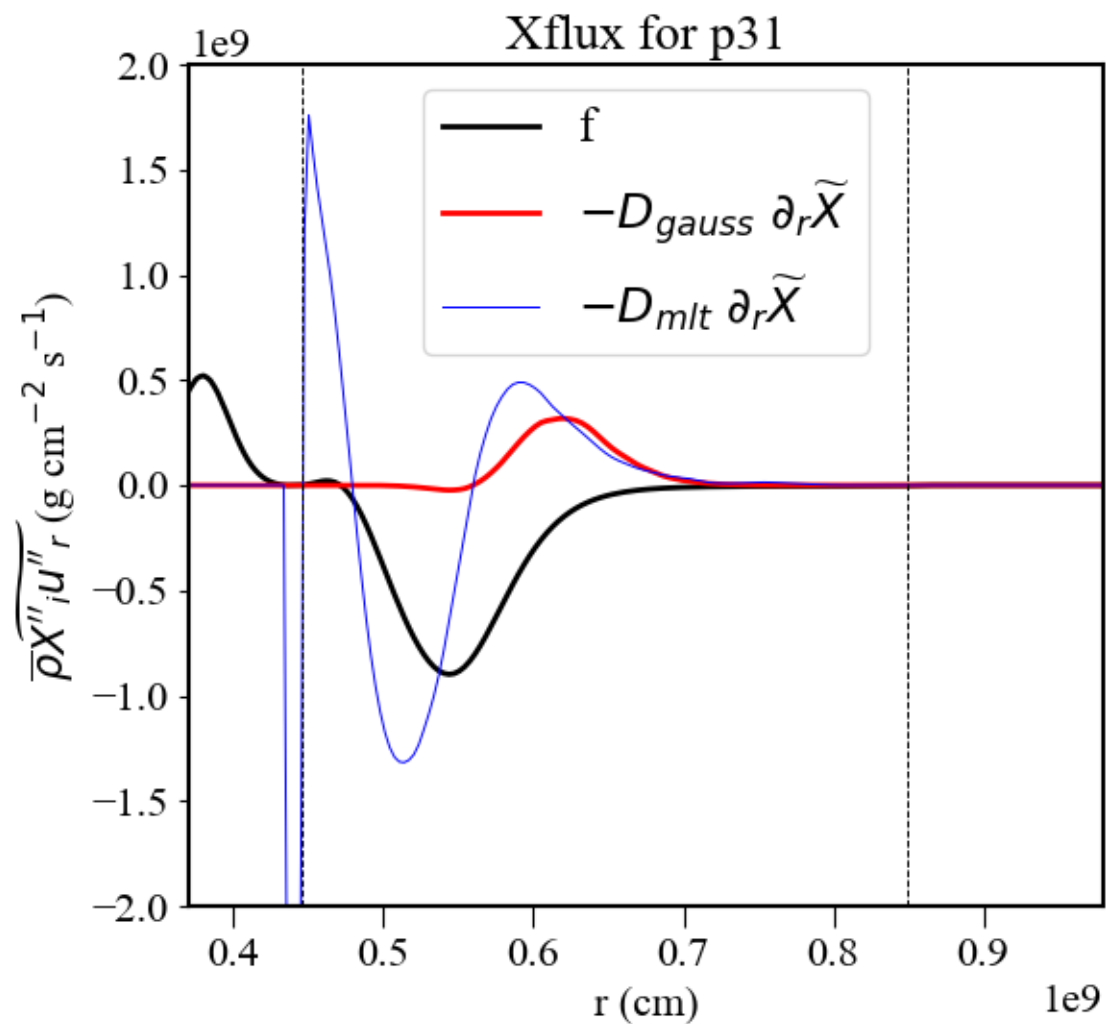
Composition flux model



Composition flux model



Composition flux model



Downgradient approximation

$$\tilde{F}_i^q \sim -\Gamma_t \frac{\partial \tilde{q}}{\partial x_i} \quad (\Gamma_t \text{ is turbulence diffusivity and } \tilde{F}_i^q = \overline{\rho q'' u_i''} \text{ is a flux of } q)$$

- can be derived from a transport equation of a diffusive passive scalar (Harlow & Hirt, 1969; Daly & Harlow, 1970):

$$\partial_t \tilde{F}_i^q - \overline{u_i'' q''} \partial_t \rho - \tilde{R}_{in} \partial_n \tilde{q} + \tilde{u}_n \overline{\rho \partial_n u_i'' q''} + \tilde{F}_n^q \partial_n \tilde{u}_i + \partial_n \overline{\rho u_n'' u_i'' q''} - \overline{u_i'' q'' \partial_n \rho u_n''} = -\overline{q''} \partial_i \bar{P} - \overline{q'' \partial_i P'} + \partial_n (\overline{\lambda \rho u_i'' \partial_n q''}) + f \tilde{F}_i^q$$

where q is the passive scalar governed by a diffusion equation $D_t q = \lambda \nabla^2 q$

It implies, that the downgradient approximation holds only for:

- a transport of a diffusive passive scalar
- a flow in steady state ($\partial_t \tilde{F}_i^q = 0$)
- an incompressible flow ($\partial_t \rho = 0$)
- a flow with no background velocities ($\tilde{u}_i = 0$)
- a flow with no pressure-scalar correlations ($\overline{q''} \partial_i \bar{P} = \overline{q'' \partial_i P'} = 0$)
- a homogeneous flow ($\partial_n \overline{\rho u_n'' u_i'' q''} = 0$)
- an isotropic flow (decay-rate assumption: $\overline{\partial_n q'' \partial_n \rho u_i''} \sim f \tilde{F}_i^q$)

But, stellar turbulent convection is:

- stratified (not homogeneous)
- anisotropic
- compressible on expanding/contracting background

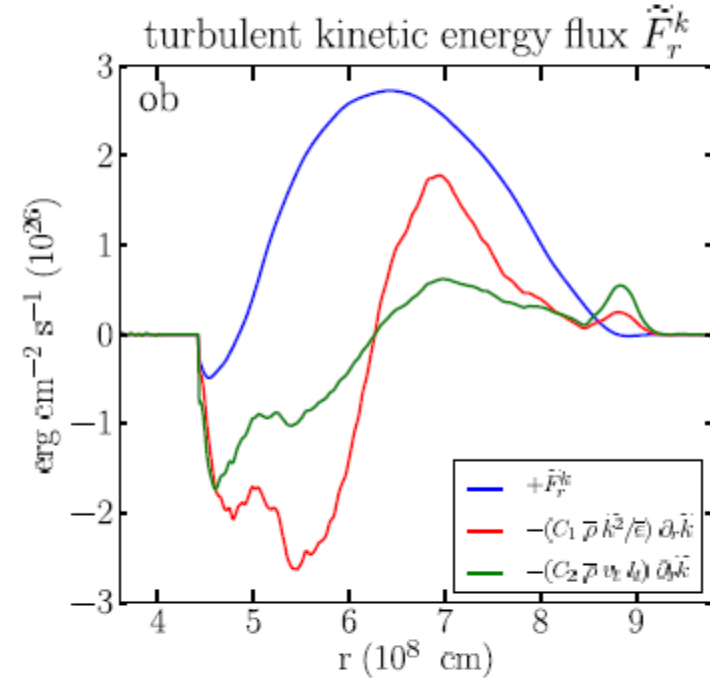


Figure 1: Downgradient approximations to the turbulent kinetic energy flux $\tilde{F}_r^k = \overline{\rho u_r'' k''}$ derived from 3D oxygen burning shell model.

- downgradient approximation is not suitable for modelling stellar processes